

Improved Theory Predictions in $b \rightarrow s\ell\ell$

Quirks in Quark Flavor Physics – 14/06/2022

Ménil Reboud

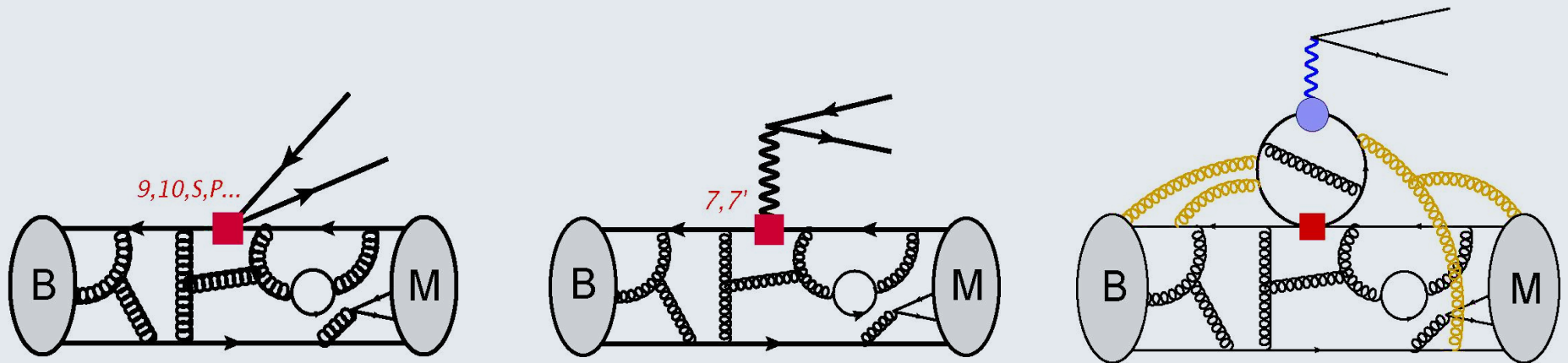
In collaboration with:

N. Gubernari, D. van Dyk, J. Virto, arXiv:2206.03797



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Form-factors in $b \rightarrow s \ell \ell$

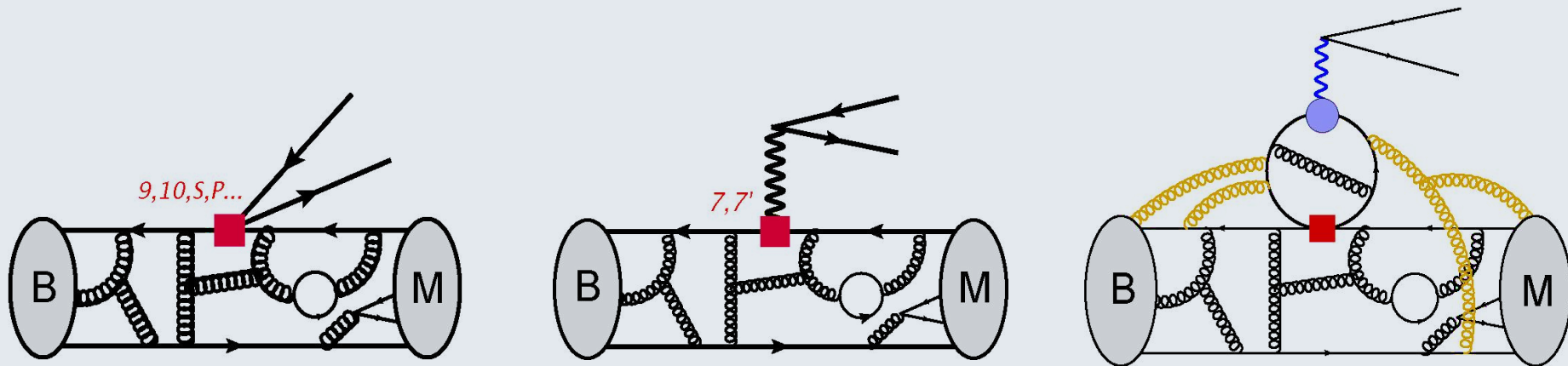


$$\mathcal{A}_\lambda^{L,R}(B \rightarrow M_\lambda \ell \ell) = \mathcal{N}_\lambda \left\{ (C_9 \mp C_{10}) \mathcal{F}_\lambda(q^2) + \frac{2m_b M_B}{q^2} \left[C_7 \mathcal{F}_\lambda^T(q^2) - 16\pi^2 \frac{M_B}{m_b} \mathcal{H}_\lambda(q^2) \right] \right\}$$

- $B \rightarrow K \mu \mu, B \rightarrow K^* \mu \mu$
- $B_s \rightarrow \varphi \mu \mu$
- $\Lambda_b \rightarrow \Lambda \mu \mu, \dots$

We focus on these **3 channels**
(and their isospin partners)!

Form-factors in $b \rightarrow s\ell\ell$



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Local form-factors

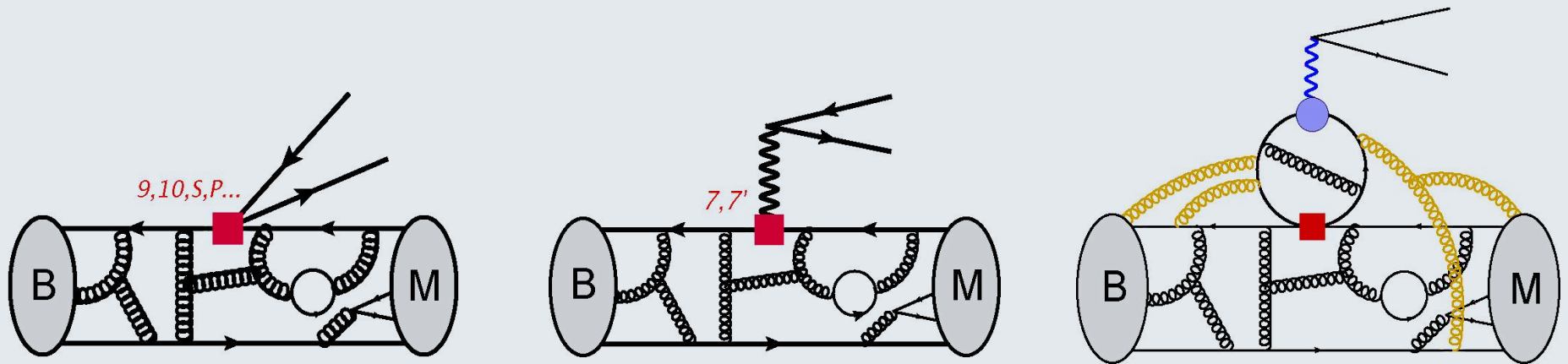
$$\mathcal{F}_{\mu}(k, q) = \langle \bar{M}(k) | \bar{s} \gamma_{\mu} b_L | \bar{B}(q+k) \rangle$$

Updated predictions based on:

- Lattice QCD calculations
- Light-cone sum rules estimates

... more in backup.

Form-factors in $b \rightarrow s\ell\ell$



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Non-local form-factors

$$\mathcal{H}_\lambda(q^2) = i\mathcal{P}_\mu^\lambda \int d^4x e^{iq \cdot x} \langle \bar{M}_\lambda(k) | T \{ \mathcal{J}_{\text{em}}^\mu(x), \mathcal{C}_i \mathcal{O}_i(0) \} | \bar{B}(q+k) \rangle$$

→ Main contributions: $\mathcal{O}_1^c, \mathcal{O}_2^c$ the so-called “charm-loops”

QCD factorization in heavy quark limit [Beneke, Feldmann, Seidel, '01 & '04]

- ◆ Uses relations between the **local** form-factors
- ◆ + **perturbative** contributions from the charm loops

→ **Limited control** on the uncertainties

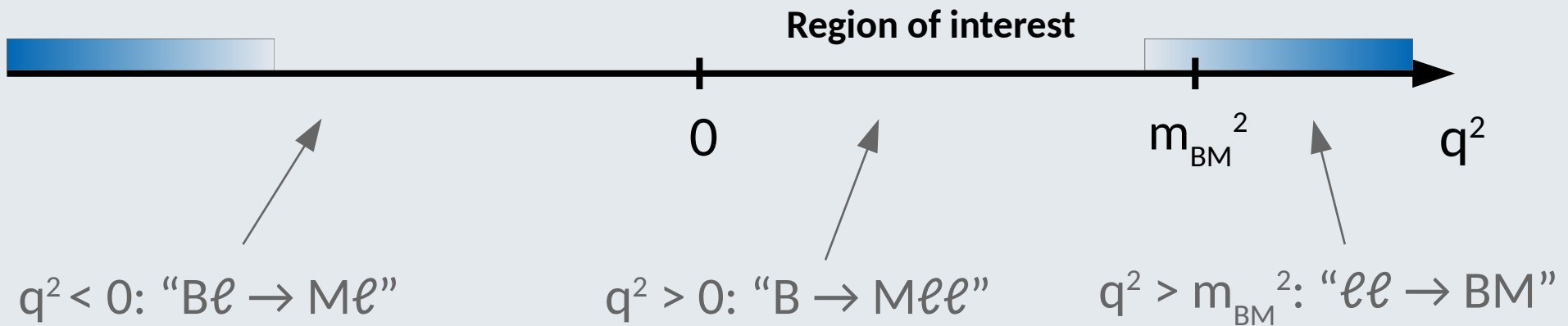
→ **Knows nothing about the J/ψ and the $\psi(2S)$!**

Constraints on H_λ

1. Two types of **OPE** can be used for H_λ :

- **Local OPE** $|q|^2 \gtrsim m_b^2$ [Grinstein, Piryol 2004][Beylich, Buchalla, Feldmann 2011]

→ We will discuss it later



Constraints on H_λ

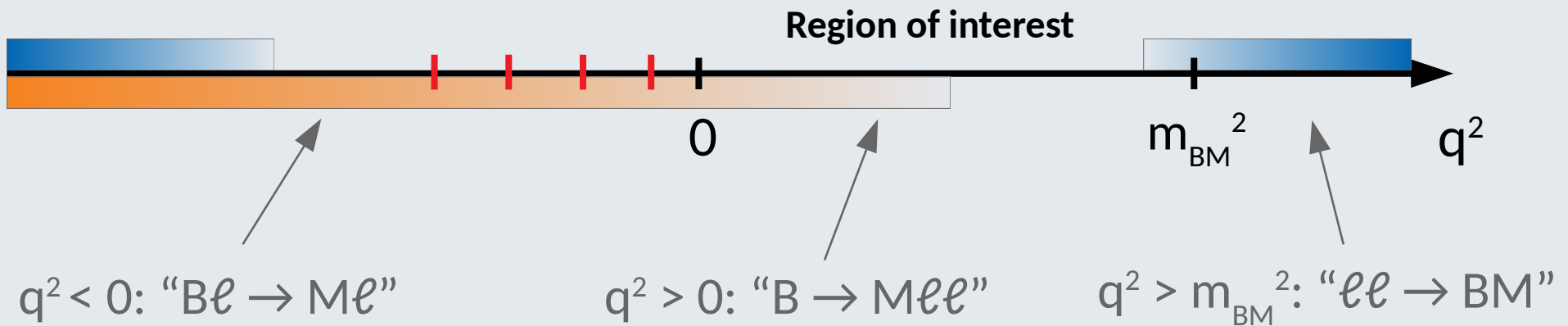
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- **Light Cone OPE** $q^2 \ll 4m_c^2$ [Khodjamirian, Mannel, Pivovarov, Wang 2010]

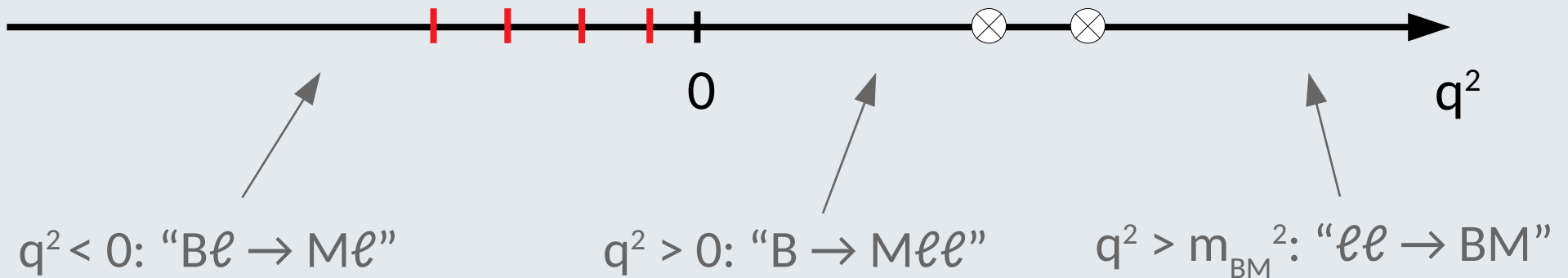
→ theory points at $q^2 < 0$ [Gubernari, van Dyk, Virto 2020]



Constraints on H_λ

2. Charmonium resonances [Bobeth, Chrzaszsz, van Dyk, Virto'17]:

- H_λ presents **poles** at $q^2 = m_{J/\psi}^2$ and $m_{\psi(2S)}^2$
- For this work we only use **$B \rightarrow M J/\psi$** data

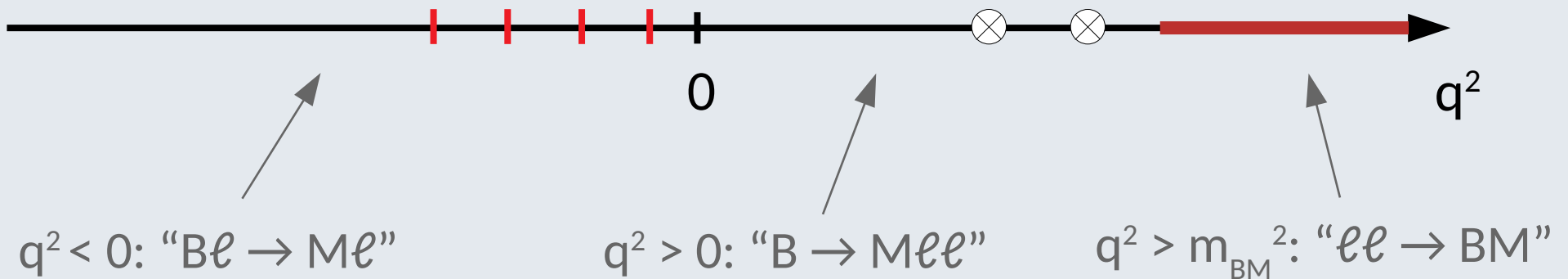


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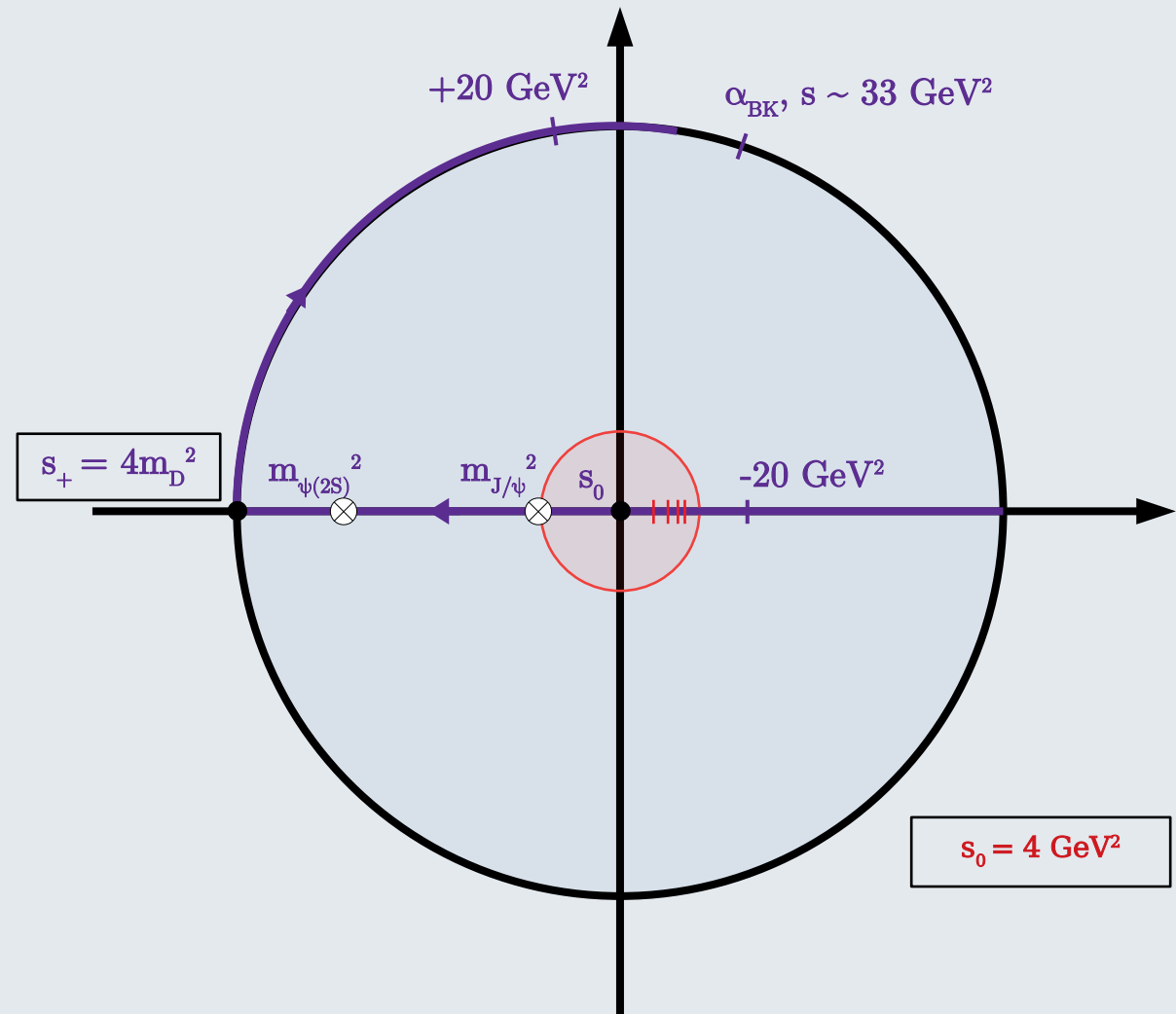
3. H_λ has a **branch cut** for $q^2 > 4m_D^2$



Parametrization of H_λ

- z-mapping

$$z(s) \equiv \frac{\sqrt{s_+ - s} - \sqrt{s_+ - s_0}}{\sqrt{s_+ - s} + \sqrt{s_+ - s_0}}$$



Parametrization of H_λ

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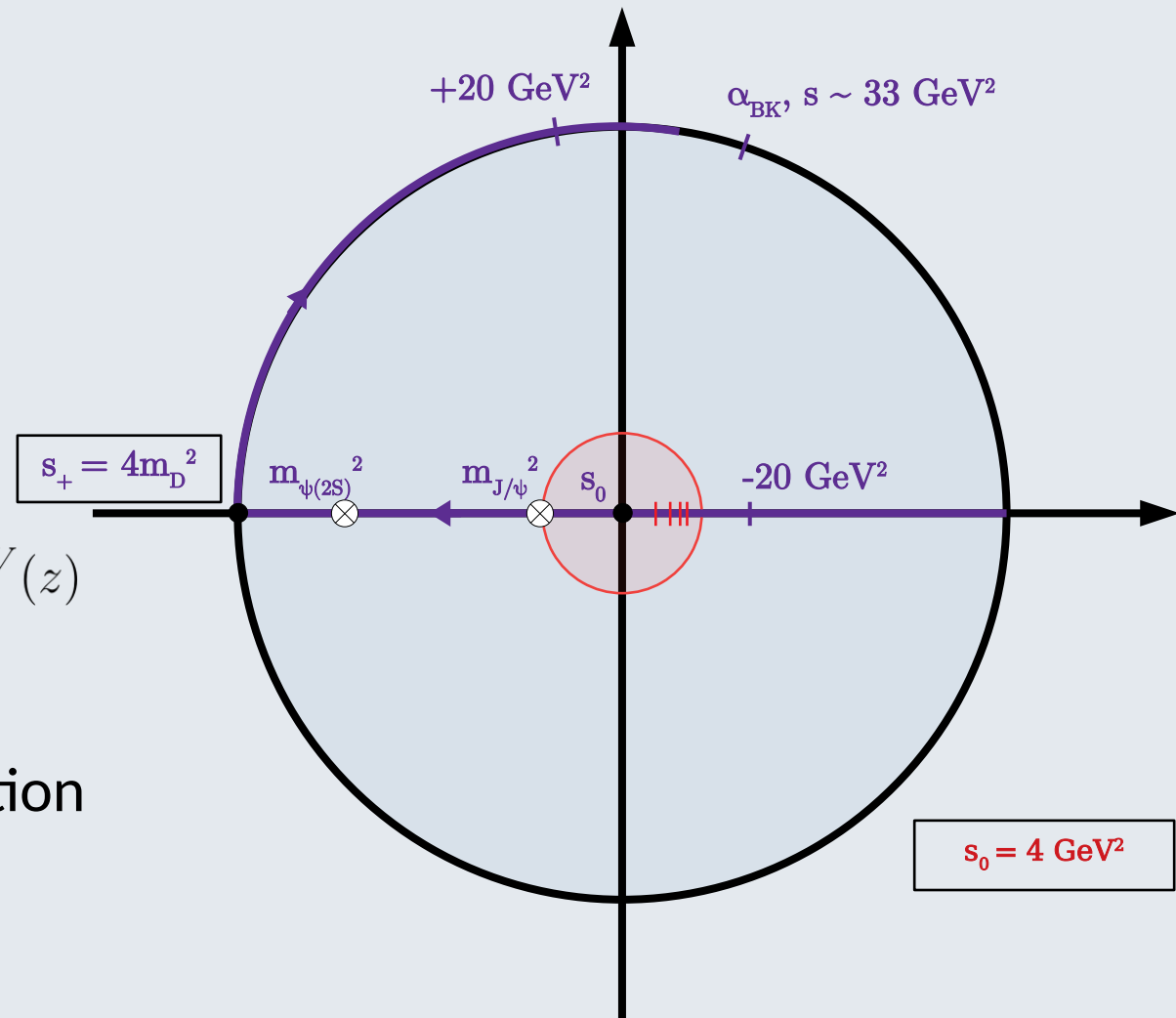
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- **Analyticity**

$$\hat{\mathcal{H}}_\lambda^{B \rightarrow V}(z) \equiv \phi_\lambda^{B \rightarrow V}(z) \mathcal{P}(z) \mathcal{H}_\lambda^{B \rightarrow V}(z)$$

→ $\mathcal{P}(z)$ captures the poles

→ $\Phi(z)$ is a useful normalization



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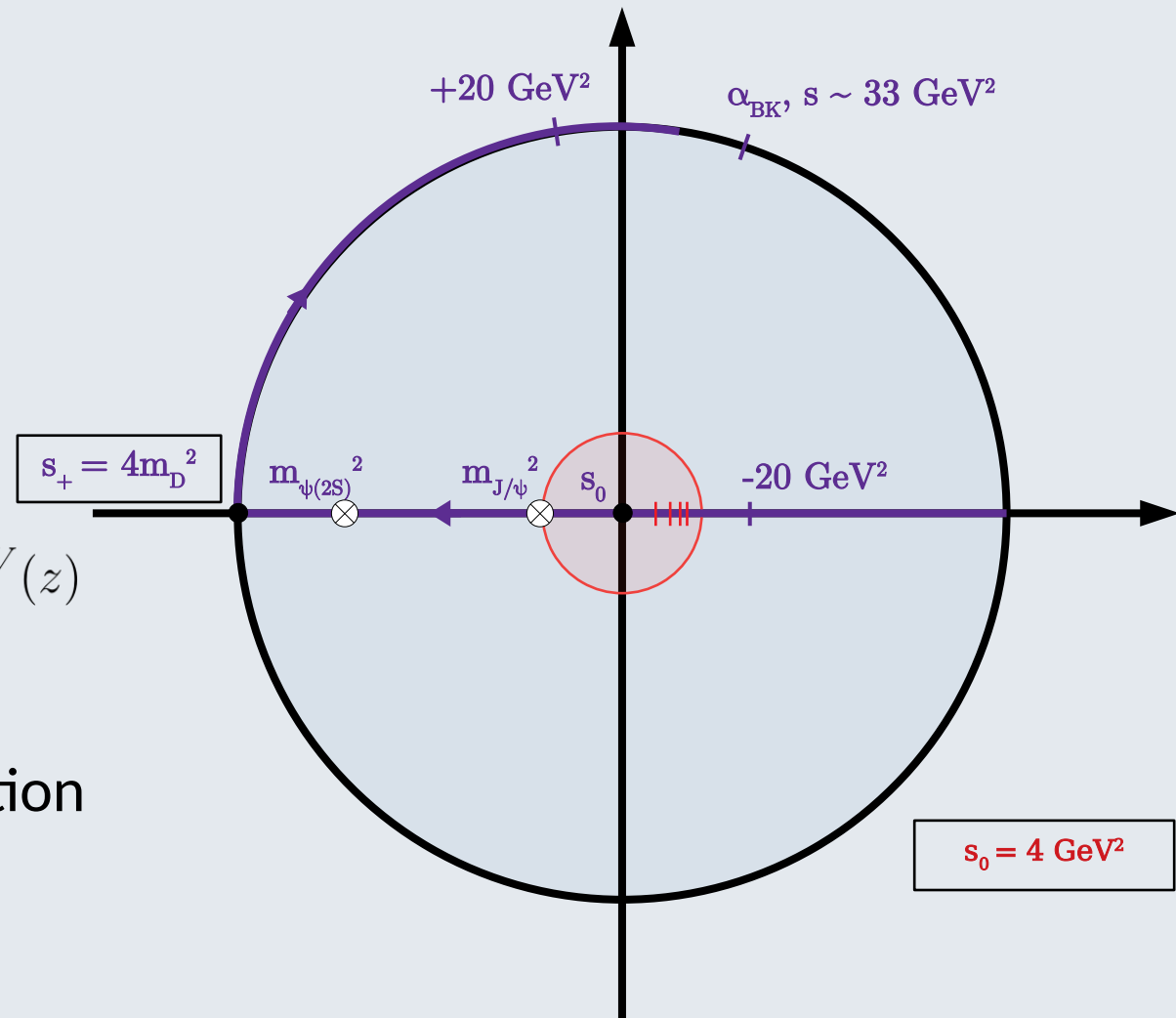
- **z-expansion**

$$\hat{\mathcal{H}}_\lambda^{B \rightarrow M}(z) = \sum_{n=0}^{\infty} a_{\lambda,n}^{B \rightarrow M} z^n$$

[Bobeth, Chrzaszcz, van Dyk, Virto 2017]

$$\hat{\mathcal{H}}_\lambda^{B \rightarrow M}(z) = \sum_{n=0}^{\infty} a_{\lambda,n}^{B \rightarrow M} p_n^{B \rightarrow M}(z)$$

[Gubernari, van Dyk, Virto, 2020]



Dispersive bound

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- **Dispersive bound** (from the **Local OPE**)

$$1 > 2 \int_{-\alpha_{BK}}^{+\alpha_{BK}} d\alpha \left| \hat{\mathcal{H}}_0^{B \rightarrow K}(e^{i\alpha}) \right|^2 + \sum_{\lambda} \left[2 \int_{-\alpha_{BK^*}}^{+\alpha_{BK^*}} d\alpha \left| \hat{\mathcal{H}}_\lambda^{B \rightarrow K^*}(e^{i\alpha}) \right|^2 + \int_{-\alpha_{B_s\phi}}^{+\alpha_{B_s\phi}} d\alpha \left| \hat{\mathcal{H}}_\lambda^{B_s \rightarrow \phi}(e^{i\alpha}) \right|^2 \right]$$

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→ With orthonormal polynomials: $\hat{\mathcal{H}}_\lambda^{B \rightarrow M}(z) = \sum_{n=0}^{\infty} a_{\lambda,n}^{B \rightarrow M} p_n^{B \rightarrow M}(z)$

$$\sum_{n=0}^{\infty} \left\{ 2 \left| a_{0,n}^{B \rightarrow K} \right|^2 + \sum_{\lambda=\perp, \parallel, 0} \left[2 \left| a_{\lambda,n}^{B \rightarrow K^*} \right|^2 + \left| a_{\lambda,n}^{B_s \rightarrow \phi} \right|^2 \right] \right\} < 1$$

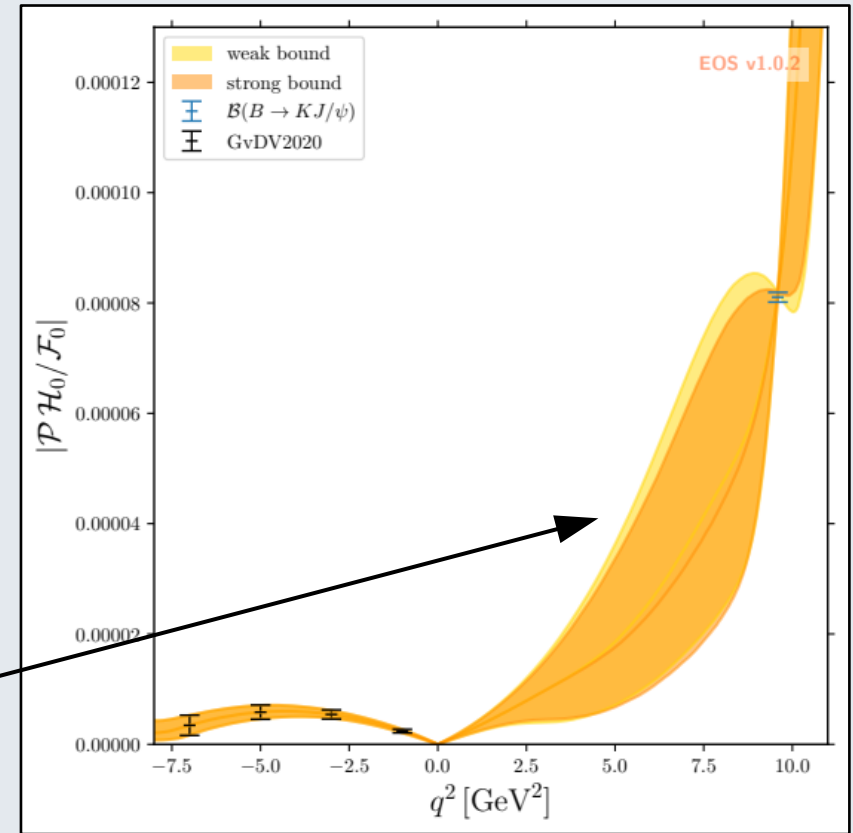
Anticipating on the results:

Expand up to **order 5**:

- 12 real parameters
- 8 constraints at negative q^2
- 1 constraint at $m_{J/\psi}^2$
→ 3 free parameters **constrained by the dispersive bound!**

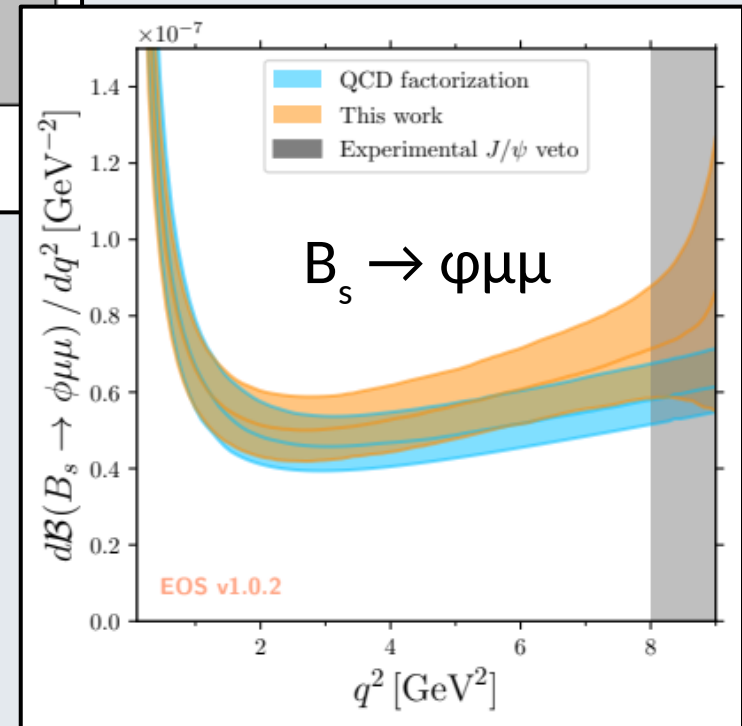
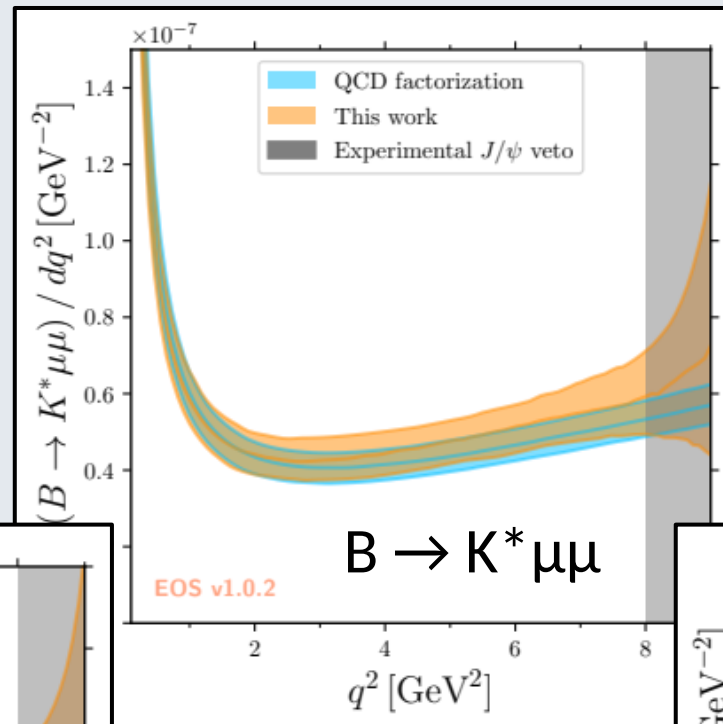
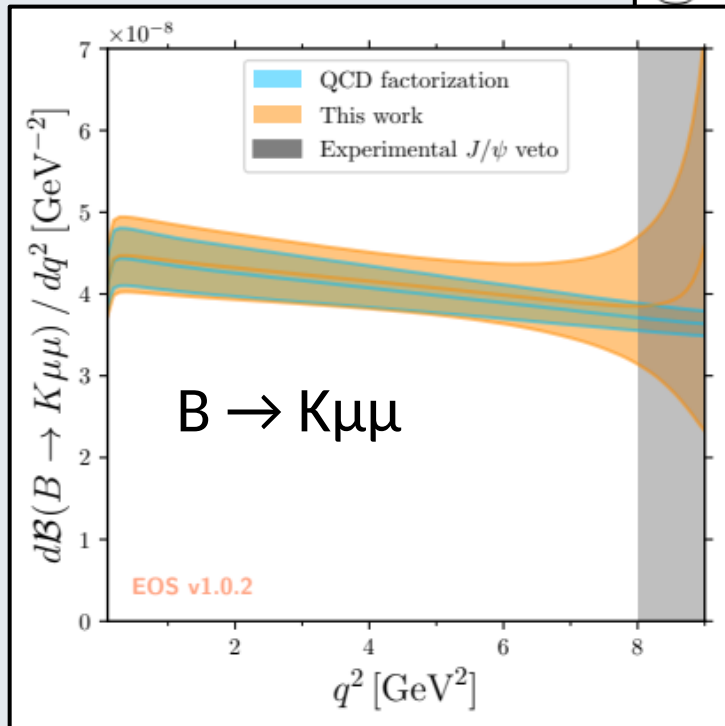
1) **Controlled uncertainty** in the physical region

2) Adding an order in the expansion **doesn't increase this uncertainty!**



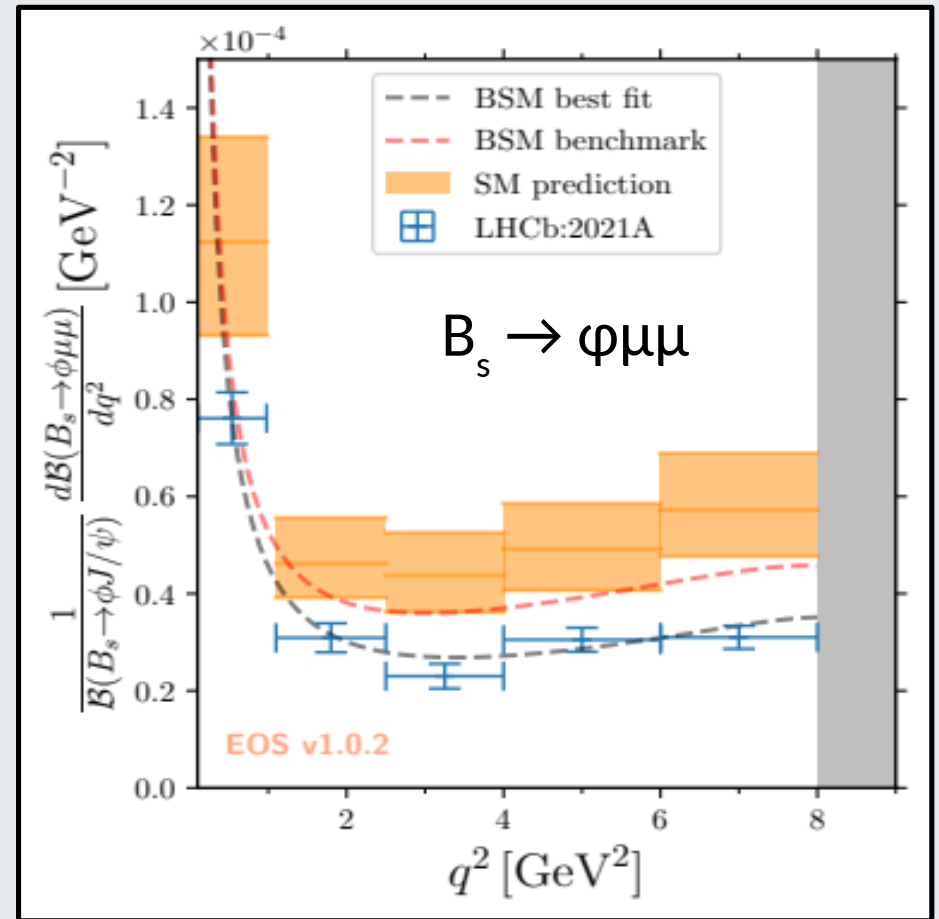
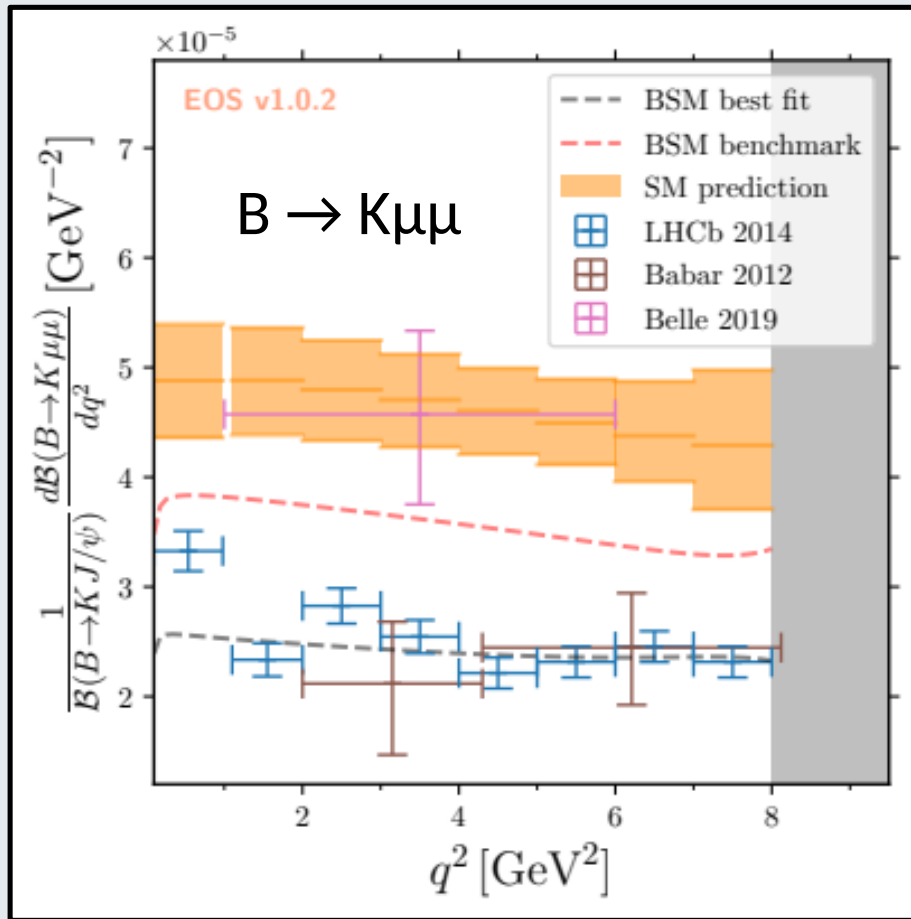
Comparison with QCDF

- Good overall agreement with QCDF
- Larger uncertainties especially near the J/ψ



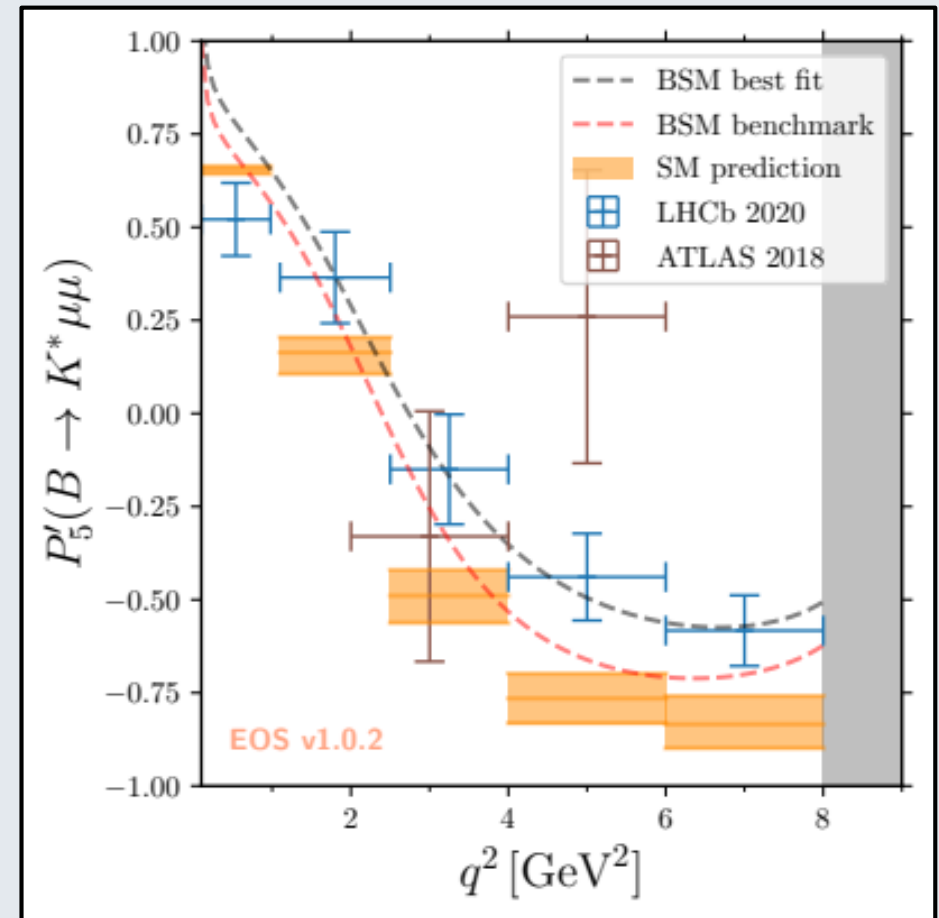
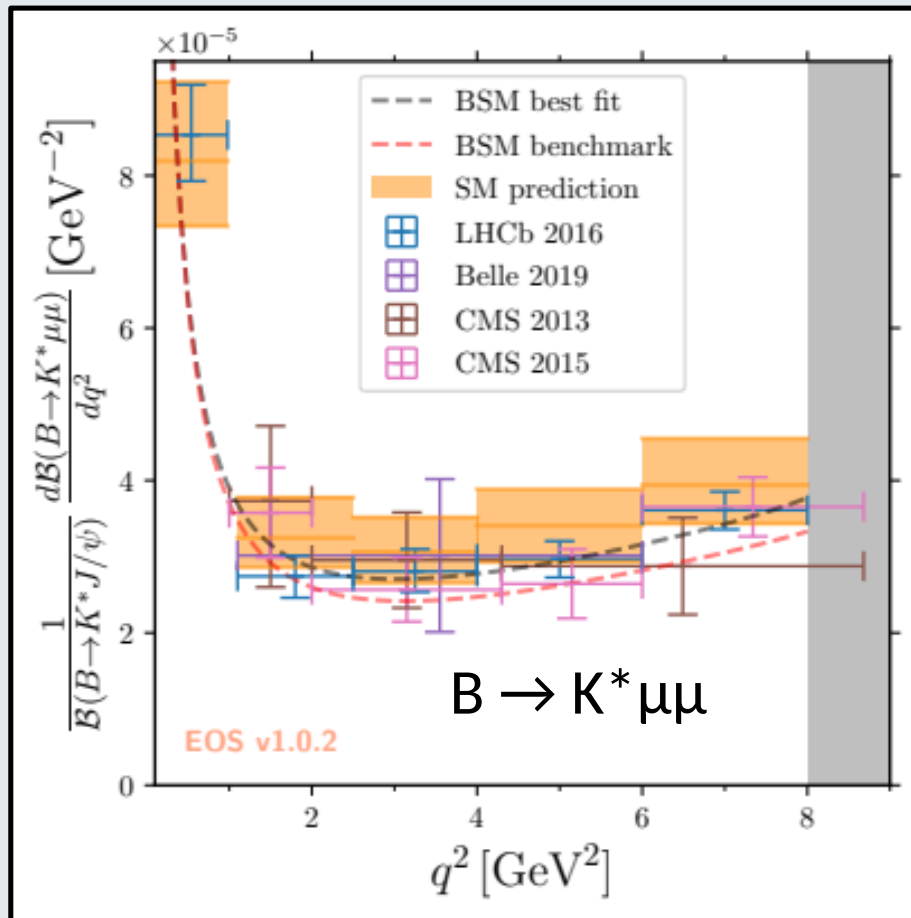
Updated (B)SM predictions

- We confirm the overall **tension** with experimental data



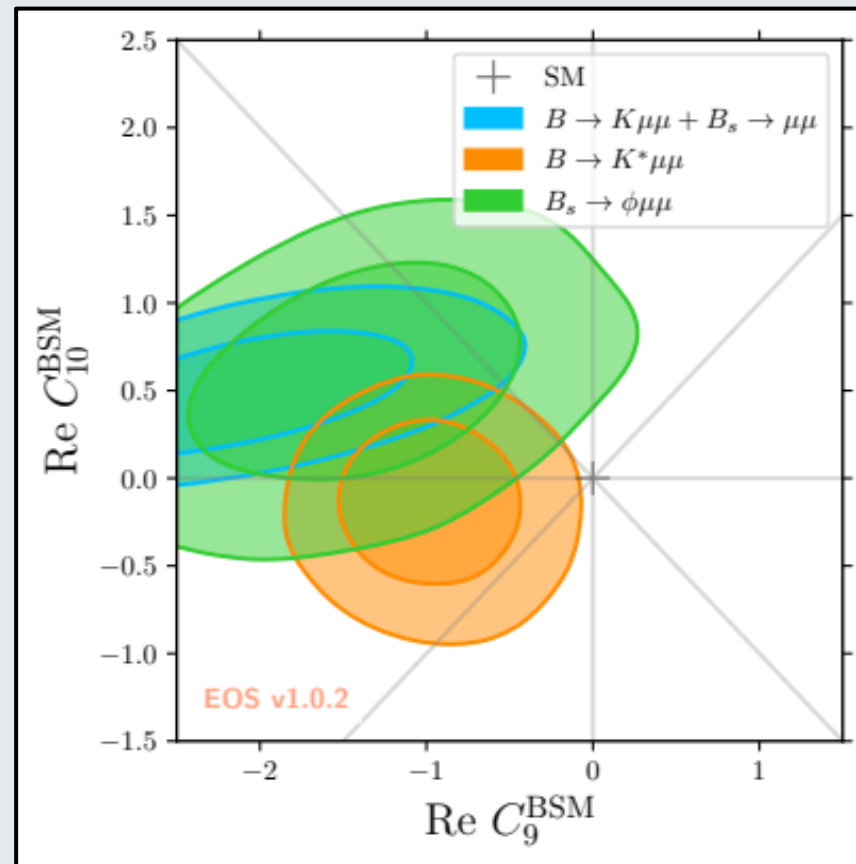
Updated (B)SM predictions

- For $B \rightarrow K^* \mu\mu$ the tension is smaller than in the literature due to **different approaches and inputs**



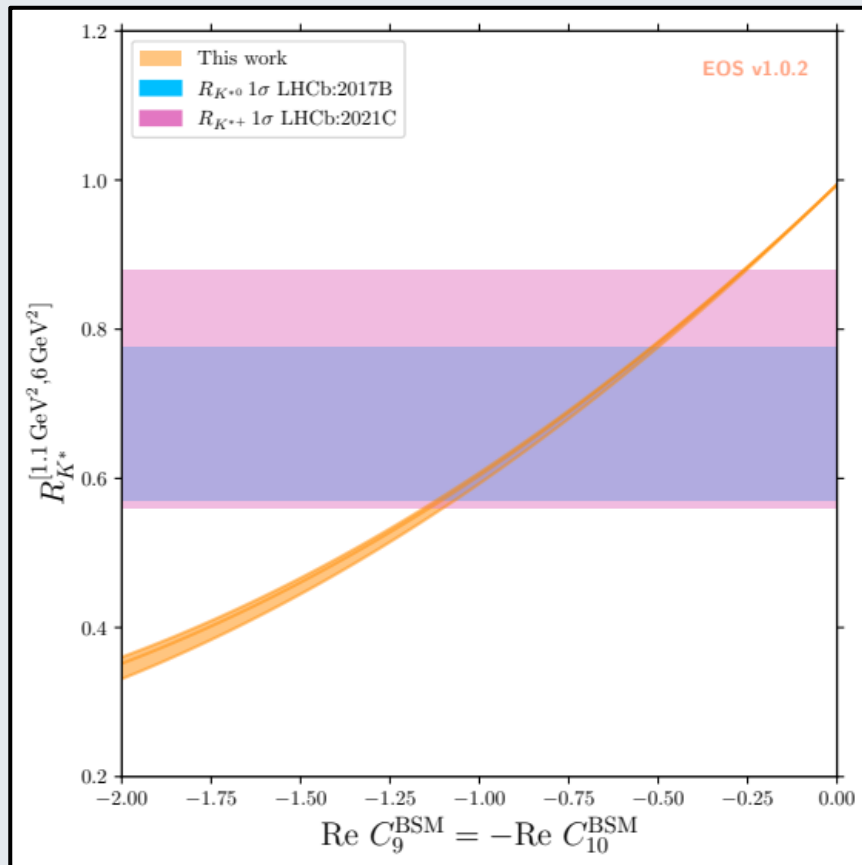
BSM analysis

- A combined BSM analysis would be **very CPU expensive** (130 nuisance parameters!)
- Fit **separately** C_9 and C_{10} for the three channels:
 $B \rightarrow K\mu^+\mu^- + B_s \rightarrow \mu^+\mu^-$, $B \rightarrow K^*\mu^+\mu^-$ and $B_s \rightarrow \phi\mu^+\mu^-$



BSM analysis

- A combined BSM analysis would be **very CPU expensive** (130 nuisance parameters!)
- Fit **separately** C_9 and C_{10} for the three channels



- Flavor universality-testing ratios

$$R_{K^{(*)}} = \frac{\mathcal{B}(B \rightarrow K^{(*)} \mu^+ \mu^-)}{\mathcal{B}(B \rightarrow K^{(*)} e^+ e^-)}$$

are **weakly sensitive to non-local contributions!**

Summary & Outlook

- We provide new **SM predictions of $b \rightarrow s\ell\ell$ observables** with
 - **Updated fit** to the *local* form-factors;
 - **Controlled uncertainties** on the *non-local* contributions.
- We performed a **(C_9, C_{10}) BSM analysis**, confirming the current trend.
- **What is next?**
 - Perform an **extended BSM analysis**;
 - Extend the prediction to **higher q^2** values (including the $\psi(2S)$)
 - Include **other channels** $\Lambda_b \rightarrow \Lambda \mu\mu, \dots$

Back-up

Putting everything together:

- The fit is performed in two steps...
 - Preliminary fits:
 - **Local** form factors:
 - BSZ parametrization (**8 + 19 + 19 parameters**)
 - LCSR + LQCD, **more in the backup**
 - **Non-local** form factors:
 - order 5 GvDV parametrization (**12 + 36 + 36 parameters**)
 - 4 points at negative $q^2 + B \rightarrow M J/\psi$ data
 - **130 nuisance parameters**
 - ‘Proof of concept’ fit to the WET’s **Wilson coefficients**
- ... using **EOS**:



EOS is a software for a variety of applications in flavour physics. It is written in C++, but provides an interface to Python.

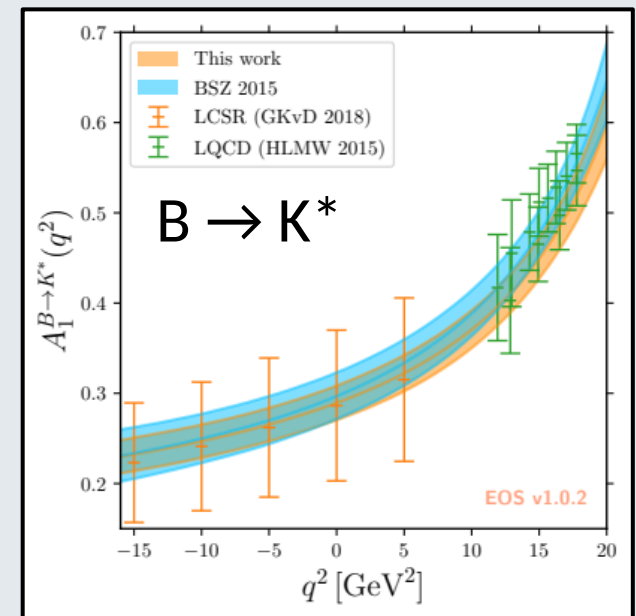
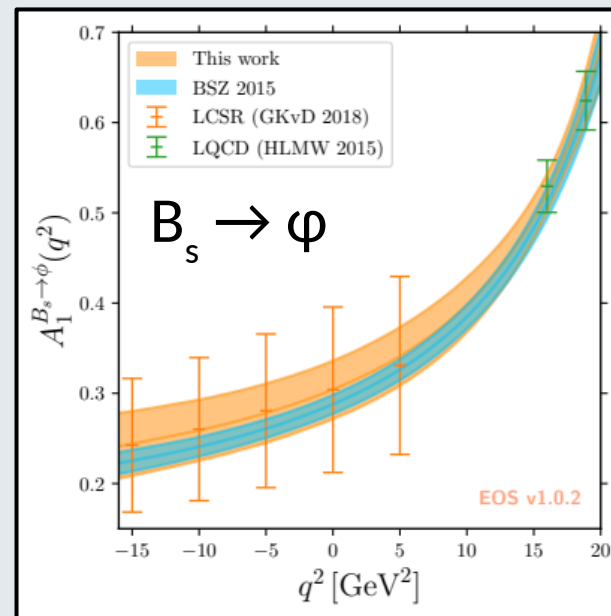
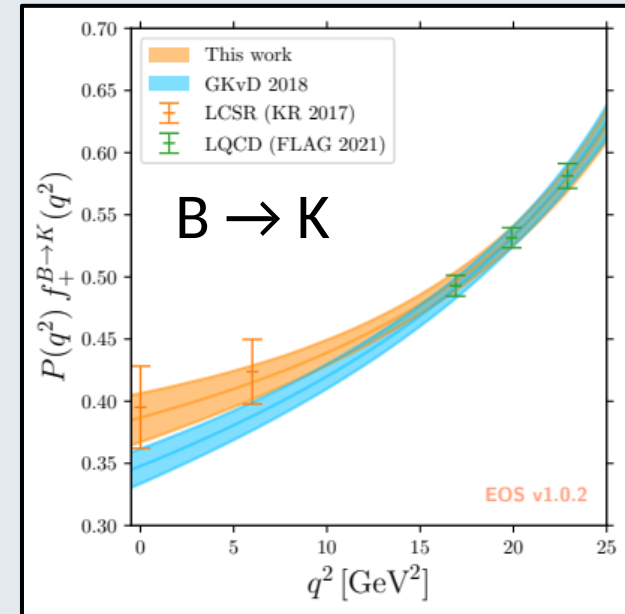


<https://eos.github.io/>

Fit to local form factors

Combined fit to LCSR and lattice:

- $B \rightarrow K$:
 - HPQCD'17; FNAL/MILC'17
 - Khodjamiriam and Rusov'17
- $B \rightarrow K^*$:
 - Horgan, Liu, Meinel and Wingate'15
 - Gubernari, Kokulu and van Dyk'18
- $B_s \rightarrow \phi$:
 - Horgan, Liu, Meinel and Wingate'15
 - Gubernari, van Dyk and Virto'20



A few remarks

1. QCD Factorization [Beneke, Feldmann, Seidel, 2001 & 2004]
2. Theory uncertainties due to charm-loops **cancel in ratios observables** → “clean” observables
Anomalies are not entirely due to charm-loops!
3. **Agreement** between “clean” and “not-so-clean” observables
Charm-loops effects cannot be very large!
4. Naively **set theory uncertainty to 0 in H_λ** :
→ Significance of the C_9 vs. C_{10} fit rises from $\sim 4\sigma$ to $\sim 8\sigma$!
This talk is not a waste of time...
5. Theory **puzzles in $b \rightarrow \bar{s}cc$** [see e.g. Lyon, Zwicky, 2014]
We need to be careful...

Data-driven (B)SM predictions

